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Majorana Fermions in Superconducting Chains

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Introduction

→ Introduction
 Finite system
 Continuous system
 Experimental realization

Quantum many-body theory:

• Fermions $|\psi(x_1, x_2)\rangle = -|\psi(x_2, x_1)\rangle$

• Bosons $|\psi(x_1, x_2)\rangle = |\psi(x_2, x_1)\rangle$

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Properties

Pauli exclusion principle

$$|\psi(x_1, x_1)\rangle = -|\psi(x_1, x_1)\rangle$$

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Properties

Pauli exclusion principle

$$|\psi(x_1, x_1)\rangle = -|\psi(x_1, x_1)\rangle$$

$$(c_i^{\dagger})^2 = c_i^2 = 0$$

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Properties

Pauli exclusion principle

$$|\psi(x_1, x_1)\rangle = -|\psi(x_1, x_1)\rangle$$

$$|\psi(x_1, x_2)\rangle = c_1^{\dagger} c_2^{\dagger} |0\rangle$$
$$|\psi(x_2, x_1)\rangle = c_2^{\dagger} c_1^{\dagger} |0\rangle$$

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Properties

Pauli exclusion principle

$$|\psi(x_1, x_1)\rangle = -|\psi(x_1, x_1)\rangle$$

$$c_1^{\dagger}c_2^{\dagger}|0\rangle = -c_2^{\dagger}c_1^{\dagger}|0\rangle$$

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Properties

Pauli exclusion principle

$$|\psi(x_1, x_1)\rangle = -|\psi(x_1, x_1)\rangle$$

$$c_1^{\dagger}c_2^{\dagger}|0\rangle = -c_2^{\dagger}c_1^{\dagger}|0\rangle$$
$$c_1^{\dagger}c_2^{\dagger} + c_2^{\dagger}c_1^{\dagger} \equiv \left\{c_1^{\dagger}, c_2^{\dagger}\right\} = 0$$

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Properties

Pauli exclusion principle

$$|\psi(x_1, x_1)\rangle = -|\psi(x_1, x_1)\rangle$$

$$\{c_i, c_j\} = \left\{c_i^{\dagger}, c_j^{\dagger}\right\} = 0 \qquad \left\{c_i, c_j^{\dagger}\right\} = \delta_{i,j}$$

Majorana fermions

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Properties

Hermitian creation operators



Chargeless \rightarrow linear combination of electron and hole

Majorana fermions

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Where to look for Majorana fermions?

Superconductivity



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Finite superconducting chain

A.Kitaev, Unpaired Majorana fermions in quantum wires, Physics-Uspekhi, 2001

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site occupation



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hopping between neighboring sites

$$\mathcal{H} = -\sum_{n}^{N} \mu c c^{\dagger} \left[-\sum_{n}^{N-1} (t c_{n}^{\dagger} c_{n+1} + h.c.) \right]$$

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addition of a cooper pair in neighboring sites

$$\mathcal{H} = -\sum_{n}^{N} \mu c c^{\dagger} - \sum_{n}^{N-1} (t c_{n}^{\dagger} c_{n+1} - \Delta c_{n} c_{n+1} + h.c.)$$

Majorana operators

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• Hermitian operators

• Majorana from different sites satisfy fermion commutation relations

• Two Majorana operators correspond to one fermion

Majorana operators

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$$\gamma_{2n-1} = \frac{c_n + c_n^{\dagger}}{2} \qquad \gamma_{2n} = i \frac{c_n - c_n^{\dagger}}{2}$$

• Hermitian operators

• Majorana from different sites satisfy fermion commutation relations

• Two Majorana operators correspond to one fermion

Finite superconducting chain

Kitaev hamiltonian

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Analytically solvable cases:

•
$$t = \Delta = 0, \mu \neq 0$$

Trivial phase $\mathcal{H} = i\mu \sum_{n=1}^{N} \gamma_{2n-1} \gamma_{2n}$

•
$$\mu = 0, t = -\Delta \neq 0$$

Topological phase $\mathcal{H} = it \sum_{n=1}^{N-1} \gamma_{2n} \gamma_{2n+1}$

Kitaev hamiltonian

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•
$$t = \Delta = 0, \mu \neq 0$$



•
$$\mu = 0, t = -\Delta \neq 0$$

Numerical calculation

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Numerical calculation

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 $t = -1, \Delta = -1, N = 20$



Infinite superconducting chain

Closing the chain

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Periodic boundary conditions

Fourier transform into momentum space

Momentum space

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$$|k
angle = N^{-1/2} \sum_{n}^{N} e^{-ikn} |n
angle$$
 Transformation

$$\mathcal{H}(k) = \sum_{k} \langle k | H | k \rangle$$

Hamiltonian becomes decomposable

$$\langle k|H|k \rangle = (-2tcos(k) - \mu)\tau_z + \Delta sin(k)\tau_y$$

Momentum space

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$$\langle k|H|k \rangle = (-2tcos(k) - \mu)\tau_z + \Delta sin(k)\tau_y$$

$$\tau_z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \qquad \tau_y = \begin{pmatrix} 0 & i \\ -i & 0 \end{pmatrix}$$

Problem becomes 2-dimentional

Energy bands

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Gap closing at μ = -2t and μ = 2t

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Quantity that can identify the topological phase of the system

$$\mathcal{H} = \bar{\sigma} \cdot \bar{h}(k) \quad \bar{\sigma} = (\tau_x, \tau_y, \tau_z)$$

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Quantity that can identify the topological phase of the system

$$\mathcal{H} = \bar{\sigma} \cdot \bar{h}(k) \quad \bar{\sigma} = (\tau_x, \tau_y, \tau_z)$$

$$\langle k|H|k \rangle = (-2t\cos(k) - \mu)\tau_z + \Delta \sin(k)\tau_y$$

 $\bar{h}(k) = (0, \ \Delta \sin(k), \ -2t\cos(k) - \mu)$

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Invariant quantity $Q = sgn(h_z(0)h_z(\pi))$

 $Q = \begin{cases} 1 & \text{In the trivial phase} \\ -1 & \text{In the topological phase} \end{cases}$

More general Hamiltonian $h_{x,y}(k) = -h(-k) \quad h_z(k) = h_z(-k)$

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 $h_x(k) = sin(3k)$ $h_y(k) = 2\Delta sin(2k)$ $h_z(k) = 2tcos(2k) - \mu$





Q = 1

Trivial phase

Experimental realization and results

s & p superconductors

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• S-pairing



• P-pairing



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Allows for momentum dependent band gap

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Spinless p-wave superconductor



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Spinless p-wave superconductor



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Superconductivity

Magnetic field

•Spin orbit coupling

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Electron chain + superconductivity



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Electron chain + superconductivity + spin orbit coupling



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Electron chain + superconductivity + spin orbit coupling + magnetic field



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Electron chain + superconductivity + spin orbit coupling + magnetic field



Condition: $B^2 > \Delta^2 + \mu^2$









Conductance Introduction Conductance Finite system Continuous system → Experimental realization



Experimental results

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V. Mourik et al., Signatures of Majorana fermions in hybrid superconductor-semiconductor nanowire devices, Science, 2012

Extensions

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•Wire circuits

- Higher dimensions
- Exchange operations

Thank you for listening